

Electrical conductivity and defect structure of lanthanum-doped (U,Pu)O_{2±x}

Toshihide Tsuji and Keiji Naito*

Department of Nuclear Engineering, Faculty of Engineering, Nagoya University, Furo-cho, Chikusa-ku, Nagoya 464-01 (Japan)

Toshiyuki Yamashita and Takeo Fujino**

Department of Chemistry and Fuel Research, Japan Atomic Energy Research Institute, Tokai-mura, Ibaraki 319-11 (Japan)

Abstract

The electrical conductivities σ of La_z(U_{0.8}Pu_{0.2})_{1-z}O_{2±x} ($z=0.01$ and 0.1) were measured in the range 10^{-15} Pa $\leq P_{O_2} \leq 10^{-1.5}$ Pa at 1273 K. Three different dependences of σ on P_{O_2} for La_z(U_{0.8}Pu_{0.2})_{1-z}O_{2±x} ($z=0.01$ and 0.1) were observed, similar to the results for the sample not containing lanthanum ($z=0$). An additional dependence of σ on P_{O_2} was observed for the sample with $z=0.1$, similar to that of the hypostoichiometric La_{0.1}U_{0.9}O_{2-x} phase, suggesting the existence of the hypostoichiometric La_{0.1}(U_{0.8}Pu_{0.2})_{0.9}O_{2-x} phase. The effect of doping on the defect structure and the electrical conduction mechanism of La_z(U_{0.8}Pu_{0.2})_{1-z}O_{2±x} ($z=0.01$ and 0.1) is discussed.

1. Introduction

Mixed oxide fuel (U,Pu)O_{2±x} is of interest in reactor engineering for use as fuel for fast breeder reactors and plutonium enriched thermal reactors. In preceding papers [1–3], we reported the electrical conductivities σ of U_{1-y}Pu_yO_{2±x} ($y=0.05, 0.10, 0.20, 0.30, 0.50$ and 0.90) under an oxygen partial pressure P_{O_2} ranging from 10^{-15} to $10^{-1.5}$ Pa at 1273 K. Three different dependences of σ on P_{O_2} were observed irrespective of y . The defect structure and the electrical conduction mechanism in these three regions were discussed.

In this study, the effect of doping on the electrical conductivities of La_z(U_{0.8}Pu_{0.2})_{1-z}O_{2±x} ($z=0.01$ and 0.1) was investigated in the oxygen partial pressure range from 10^{-15} to $10^{-1.5}$ Pa at 1273 K.

2. Experimental details

Purified uranium, plutonium and lanthanum nitrate solution were mixed in appropriate concentrations, and the mixed solution was gently evaporated and dried in a mantle heater. Powder of the mixed oxide was obtained by calcining the solid nitrate at 1073 K in air. After pressing the powder at about 20 kg mm⁻², the pellet was sintered in vacuum at 1673 K for 3–4 h. The

electric conductivity was measured and the oxygen partial pressure was controlled in the range 10^{-15} – $10^{-1.5}$ Pa as described in our previous papers [1–3].

3. Results and discussion

The dependences on oxygen partial pressure of the electrical conductivity for La_z(U_{0.8}Pu_{0.2})_{1-z}O_{2±x} ($z=0.01$ and 0.1) samples obtained at 1273 K are shown in Fig. 1, where the result of $\log \sigma$ vs. $\log P_{O_2}$ for undoped U_{0.8}Pu_{0.2}O_{2±x} obtained by the present authors [2] is shown for comparison. In Fig. 1, three different dependences of σ on P_{O_2} for lanthanum-doped samples ($z=0.01$ and 0.10) are seen, similar to that for the undoped sample ($z=0$). (1) Above $P_{O_2}=10^{-2}$ Pa (region I), the slope of the $\log \sigma$ vs. $\log P_{O_2}$ curve decreases with increasing oxygen partial pressure, presumably owing to the onset of the phase transition from a single-phase MO_{2±x} to a two-phase MO_{2±x}–M₄O_{9-x} region. (2) Between $P_{O_2}=10^{-4.5}$ and 10^{-2} Pa (region II), $\log \sigma$ increases linearly with increasing $\log P_{O_2}$, and the slope is smaller for the specimens with larger z values. (3) Between $P_{O_2}=10^{-5}$ and 10^{-11} Pa (region III), the electrical conductivity is independent of P_{O_2} . An additional dependence of σ on P_{O_2} is observed for the sample with $z=0.1$ below $10^{-11.5}$ Pa (region IV) as seen in Fig. 1, where the electrical conductivity increases with increasing oxygen partial pressure. This behavior is similar to that of the hypostoichiometric La_{0.1}U_{0.9}O_{2-x}

*Emeritus Professor of Nagoya University.

**Present address: Institute for Advanced Materials Processing, Tohoku University, 2-1-1 Katahira, Sendai 980, Japan.

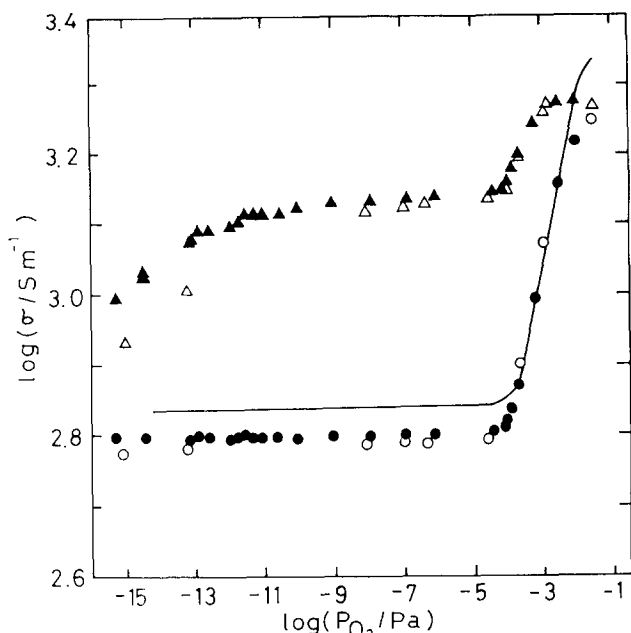
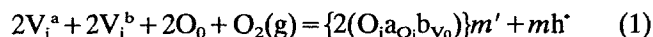


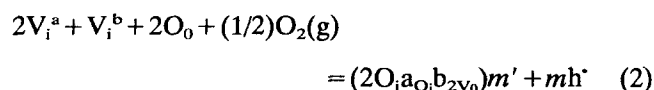
Fig. 1. Dependence on oxygen partial pressure of the electrical conductivity of La_x(U_{0.8}Pu_{0.2})_{1-x}O_{2±x} for z=0.01 (○, ●) and 0.1 (△, ▲) at 1273 K. The result of log σ vs. log P_{O2} for undoped U_{0.8}Pu_{0.2}O_{2±x} (—) obtained by us [2] is shown for comparison. The electrical conductivities were reversible for increasing (open symbols) and decreasing (closed symbols) oxygen partial pressures, except for the low oxygen partial pressure region for the sample with z=0.10.

phase [4], suggesting the existence of the hypostoichiometric La_{0.1}(U_{0.8}Pu_{0.2})_{0.9}O_{2-x} phase.

From the slope in region II shown in Fig. 1, the values of n in the relation σ ∝ P_{O2}^{1/n} for z=0, 0.01 and 0.10 were calculated as 4.1, 5.1 and 9.9 respectively. The formation of the complex defect {2(O_ia_Ob_{vO})}m', similar to the case for undoped (U,Pu)O_{2+x} [1-3], is represented by



where h* is a hole and m is the charge of the complex defect. The dependences of n=4.1 (for z=0) and n=5.1 (for z=0.01), where n=m+1, observed for La_x(U_{0.8}Pu_{0.2})_{1-x}O_{2+x} samples from the log σ-log P_{O2} relation can be found from eqn. (1) by taking the values m=3 and m=4 respectively. For the defect structure of the sample with z=0.1 where the deviation from the stoichiometric composition (x=0) is considered to be small, the formation of the complex defect (2O_ia_Ob_{2vO})m' is assumed:



The dependences of σ on P_{O2}, n=10, where n=m+1, can be found from eqn. (2) taking m=9.

In regions III and IV, the dependence of oxygen partial pressure on the compositional deviation x is needed in order to discuss the defect structure.

The electrical conductivities at 1273 K around P_{O2}=10⁻⁹ Pa which are independent of P_{O2} in Fig. 1 are shown in Fig. 2, where the solid line represents the theoretical change in electrical conductivity with the atomic ratio [M]/[M+U] (M≡Pu) reported by us [2]. The theoretical line is calculated using the disproportionation reaction Pu⁴⁺ + U⁴⁺ = Pu³⁺ + U⁵⁺. In the figure, the electrical conductivities of UO_{2±x} doped with lanthanum [4], gadolinium [5] and yttrium [6] at 1273 K are also shown for comparison. As seen in Fig. 2, the electrical conductivities of La_x(U_{0.8}Pu_{0.2})_{1-x}O_{2±x} (z=0.1) and yttrium-doped UO_{2±x} deviate from the theoretical curve calculated by the disproportionation reaction.

According to a hopping model in UO_{2±x} proposed by Aronson et al. [7], when an interstitial oxygen ion is introduced in UO₂, two U⁵⁺ ions are formed, and each U⁵⁺ ion is considered as a site for one hole which can jump to a U⁴⁺ site. If we apply this model to hyperstoichiometric (U_{1-y}Pu_y)O_{2+x} by assuming the plutonium to be tetravalent, the following equation can be derived:

$$\sigma T = \sigma_0(2x)(1-y-2x) \exp(-E_a/kT) \quad (3)$$

where E_a is the activation energy for the hopping of holes. In region III, if x=0 is assumed as a crude

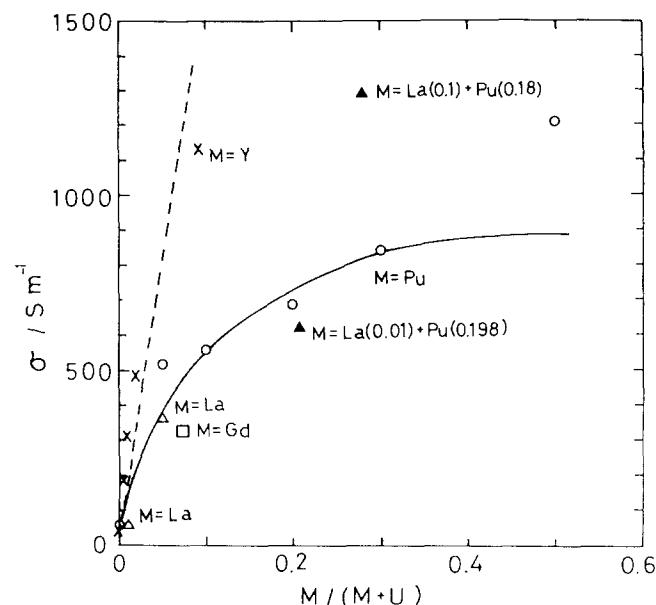


Fig. 2. Electrical conductivities at 1273 K around P_{O2}=10⁻⁹ Pa in Fig. 1 as a function of the atomic ratio [M]/[M+U]. In the figure the solid and dotted lines represent the theoretical changes in electrical conductivity calculated by the disproportionation reaction [2] and by the hopping model for holes between U⁵⁺ and U⁴⁺ ions respectively: ○, M≡Pu, ▲, M≡La + Pu, △, M≡La, □, M≡Gd, ×, M≡Y.

approximation, the ionic species of rare earth doped uranium dioxide $R_zU_{1-z}O_2$ can be expressed as $R_z^{3+}U_{1-2z}^{4+}U_z^{5+}O_2^{2-}$. If we use eqn. (3) with $E_a=0.3$ eV and $\sigma_0=3.8\times 10^8$ S m⁻¹ for UO_{2+x} from Aronson et al. [7] and substitute z and $1-2z$ for $2x$ and $1-y-2x$ respectively, then a theoretical line for the electrical conductivity for the hopping of holes between U⁵⁺ and U⁴⁺ ions can be calculated as shown by the dotted line in Fig. 2. The theoretical line fits well the results for yttrium-doped UO₂ and is larger than the experimental values for gadolinium- and lanthanum-doped UO₂. A similar calculation for La_{0.1}³⁺U_{0.62}⁴⁺U_{0.1}⁵⁺-Pu_{0.18}⁴⁺O₂²⁻ yields the calculated electrical conductivity 1.20×10^3 S m⁻¹, which is certainly comparable with the observed value of 1.29×10^3 S m⁻¹. Therefore, in yttrium-doped UO₂ and lanthanum-doped (U,Pu)O₂, electrical conduction is considered to be caused mainly by the hopping of holes between U⁴⁺ and U⁵⁺ ions

resulting from the charge neutrality condition, and the contribution from the disproportionation reaction is small.

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